

Logistics

- We have assigned Grad students to papers and are now figuring out the dates.
- Include your team ID in your github readme.
- Project 1: Papers due on February 10th.
- February 3rd is the drop without a 'W' grade and full refund deadline.

Logistics

- Student Presentation on Wednesday: Lorenz, E., “Computational Chaos”, *Physica D*, 1988
- I will post the presentation review forms this afternoon.

Lecture 5

Fixed Point Analysis

Basic Concepts: Higher Order Equations

Consider a system defined by a higher-order differential equation such as

$$\frac{d^3 x}{dt^3} = f(x, \dot{x}, \ddot{x})$$

Basic Concepts: Higher Order Equations

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We can always turn this into a system of 1st-order differential equations

Just by renaming the derivative terms:

Basic Concepts: Higher Order Equations

Consider a system defined by a higher-order differential equation such as

$$\frac{d^3 x}{dt^3} = f(x, \dot{x}, \ddot{x})$$

Substitute

We can always turn this into a system of 1st-order differential equations

Just by renaming the derivative

terms:

$$x_1(t) = x(t)$$

$$x_2(t) = \dot{x}(t)$$

$$x_3(t) = \ddot{x}(t)$$

Basic Concepts: Higher Order Equations

Consider a system defined by a higher-order differential equation such as

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$x_1(t) = x(t)$ Rewrite as a system

$x_2(t) = \dot{x}(t)$ of equations

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$$x_1(t) = x(t) \quad \text{Rewrite as a system}$$

$$x_2(t) = \dot{x}(t) \quad \text{of equations}$$

$$x_3(t) = \ddot{x}(t)$$

$$\frac{dx_1}{dt} = x_2$$

$$\frac{dx_2}{dt} = x_3$$

$$\frac{dx_3}{dt} = f(x_1, x_2, x_3)$$

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$$\frac{dx_3}{dt} = f(x_1, x_2, x_3)$$

Basic Concepts: Higher Order Equations

This is generally true for
differential equations of
any order:

so let's use
a more general
and compact notation.

Basic Concepts: Higher Order Equations

$$\frac{d\vec{x}}{dt} = \frac{d}{dt} \begin{bmatrix} x_1 \\ x_2 \\ \vdots \\ x_m \end{bmatrix} = \begin{bmatrix} x_2 \\ x_3 \\ \vdots \\ f(x_1, x_2, \dots, x_m) \end{bmatrix}$$

so let's use
a more general
and compact notation.

Basic Concepts: Higher Order Equations

$$\begin{bmatrix} x_1(t+1) \\ x_2(t+1) \\ \vdots \\ x_m(t+1) \end{bmatrix} = \begin{bmatrix} x_2(t) \\ x_3(t) \\ \vdots \\ f(x_1(t), x_2(t), \dots, x_m(t)) \end{bmatrix}$$

or in the case of a map.

Basic Concepts: Phase Space

Definition: The space spanned by all allowed values of $x_1 \dots x_m$ in a system defined by:

$$\frac{d\vec{x}}{dt} = \frac{d}{dt} \begin{bmatrix} x_1 \\ x_2 \\ \vdots \\ x_m \end{bmatrix} = \begin{bmatrix} x_2 \\ x_3 \\ \vdots \\ f(x_1, x_2, \dots, x_m) \end{bmatrix}$$

Basic Concepts: Trajectory or Orbit

Definition: A solution in the system defined by:

$$\frac{d\vec{x}}{dt} = \frac{d}{dt} \begin{bmatrix} x_1 \\ x_2 \\ \vdots \\ x_m \end{bmatrix} = \begin{bmatrix} x_2 \\ x_3 \\ \vdots \\ f(x_1, x_2, \dots, x_m) \end{bmatrix}$$

and initial conditions: $x_1 = c_1, \dots, x_m = c_m$.

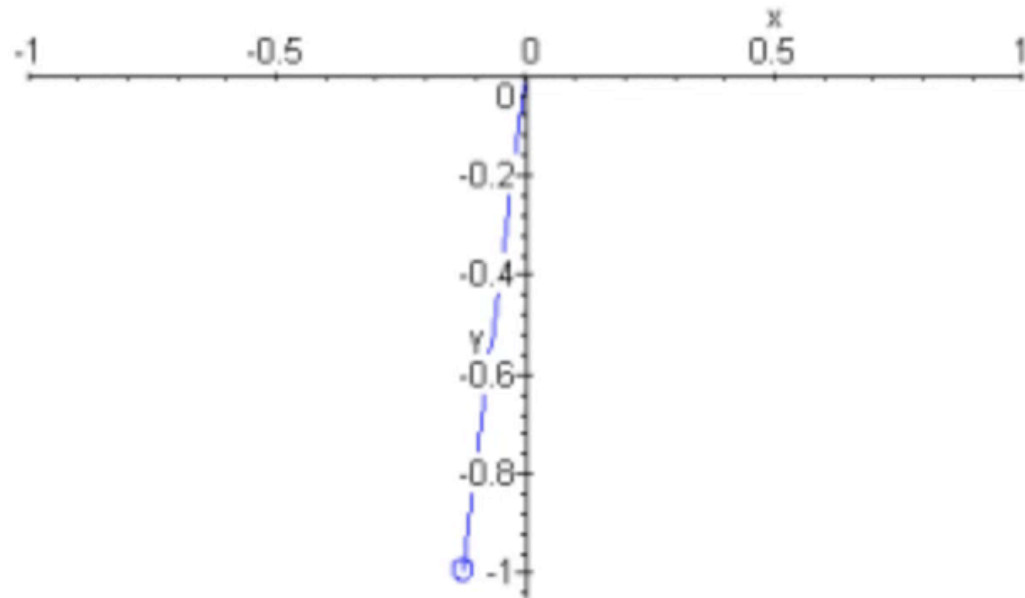
Basic Concepts: Recall from last time

System of equations for an undamped pendulum:

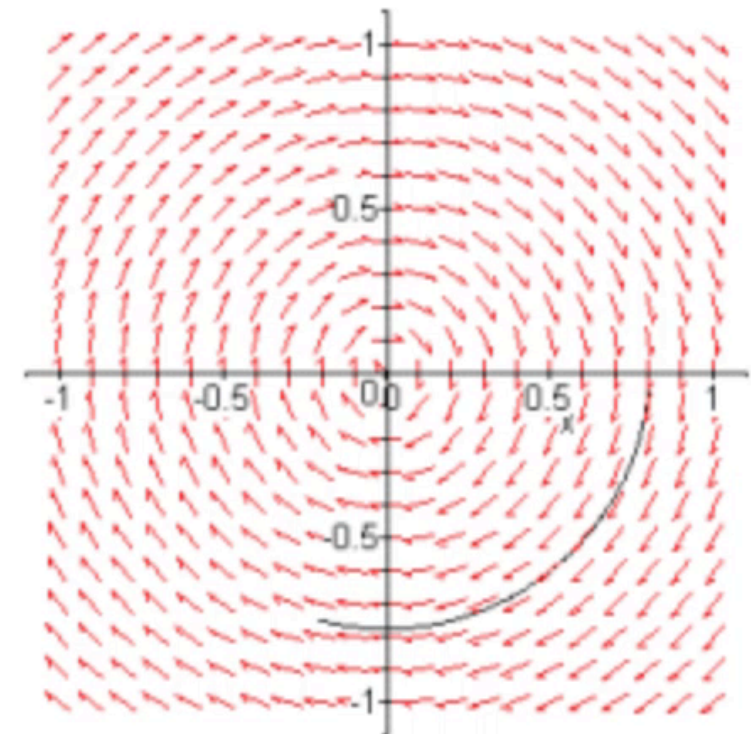
$$y_1' = y_2$$

$$y_2' = -\sin(y_1)$$

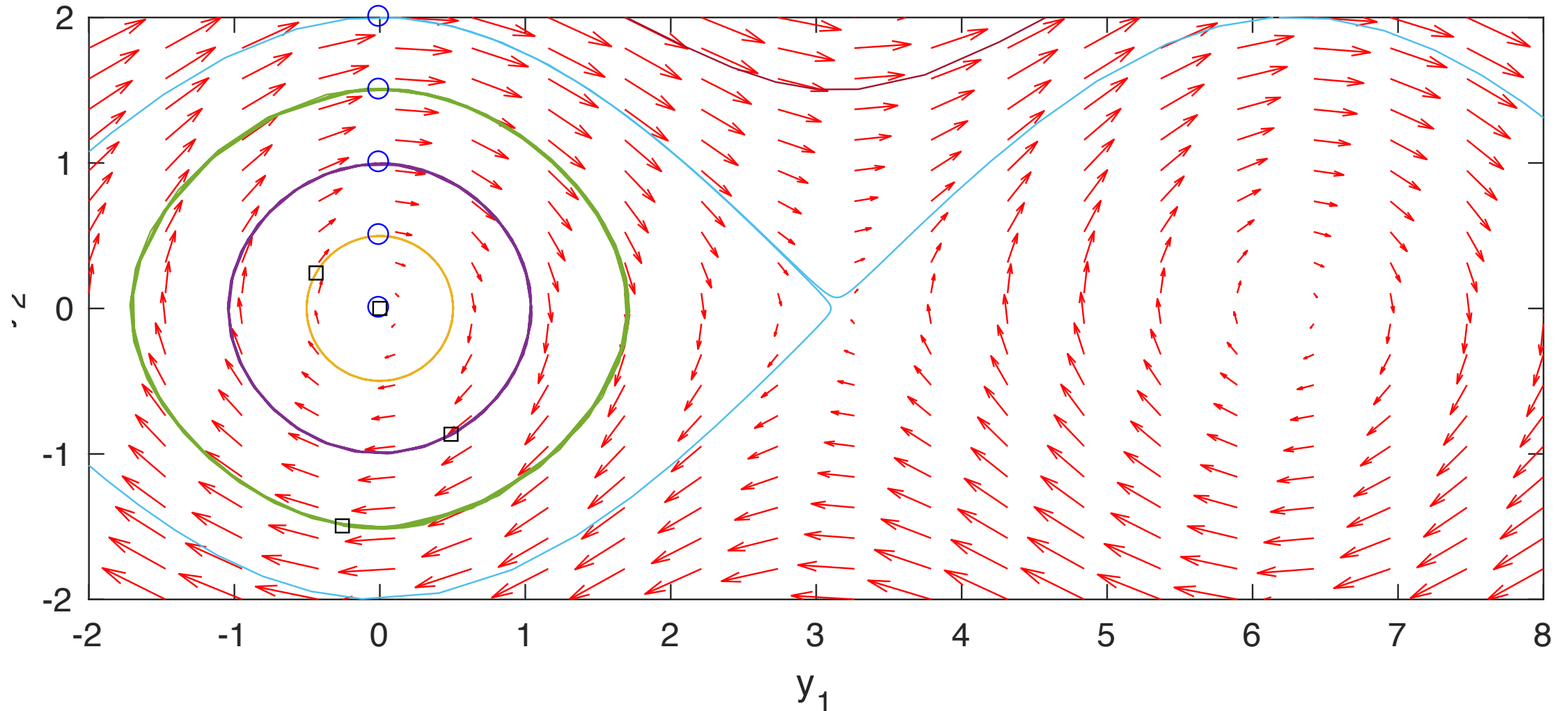
Simple Pendulum in Real Space : Undamped



Simple Pendulum in Phase Space : Undamped



Basic Concepts: What do the three fixed points mean physically?



Basic Concepts: Higher Order Equations

Since we can rewrite any higher-order differential equation as a system of 1st-order equations we only have to consider 1st-order equations from now on.

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How does this relate to part 1 of the project?

Fixed Points and Stability

Consider a fixed point x^* for a 1D system:

$$\dot{x} = f(x)$$

$$f(x^*) = 0$$

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The stability of x^* depends on the direction of the flow nearby.

Fixed Points and Stability

Why?

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Fixed Points and Stability

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To find the direction near x^*
we linearise \dot{x} near x^*

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Recall from Calculus I the Taylor expansion:

$$\frac{d}{dt}(x - x^*) = \dot{x} \approx f(x^*) + f'(x^*)(x - x^*) + \dots$$

Fixed Points: Classification

Our system: $\dot{x} = f(x)$, $f(x^*) = 0$

Plugging in and neglecting higher order terms:

$$\frac{d}{dt}(x - x^*) = \dot{x} \approx f(x^*) + f'(x^*)(x - x^*) + \dots$$

$$\frac{d}{dt}(x - x^*) \approx 0 + f'(x^*)(x - x^*)$$

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$$\frac{d}{dt}(x - x^*) = \dot{x} \approx f(x^*) + f'(x^*)(x - x^*) + \dots$$

$$\frac{d}{dt}(x - x^*) \approx 0 + f'(x^*)(x - x^*)$$

Fixed Points: Classification

Solving: $\frac{d}{dt}(x - x^*) \approx f'(x^*)(x - x^*)$

Relabel: $\Delta x := (x - x^*)$

Solve: $\frac{d\Delta x}{dt} \approx f'(x^*)\Delta x$

Fixed Points: Classification

Solve: $\frac{d\Delta x}{dt} \approx f'(x^*)\Delta x$

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$$d\Delta x \approx f'(x^*)\Delta x dt$$

Fixed Points: Classification

Solve: $\frac{d\Delta x}{dt} \approx f'(x^*)\Delta x$

$$d\Delta x \approx f'(x^*)\Delta x dt$$

$$\frac{d\Delta x}{\Delta x} \approx f'(x^*)dt$$

Fixed Points: Classification

Solve: $\frac{d\Delta x}{dt} \approx f'(x^*)\Delta x$

$$d\Delta x \approx f'(x^*)\Delta x dt$$

$$\frac{d\Delta x}{\Delta x} \approx f'(x^*) dt$$

$$\int \frac{d\Delta x}{\Delta x} \approx \int f'(x^*) dt$$

Fixed Points: Classification

Solve: $\frac{d\Delta x}{dt} \approx f'(x^*)\Delta x$ $\int \frac{1}{\Delta x} d\Delta x \approx \int f'(x^*) dt$

$$d\Delta x \approx f'(x^*)\Delta x dt$$

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$\Delta x \approx e^{f'(x^*) t}$

Fixed Points: Classification

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$$\Delta x \approx e^{f'(x^*)t}$$

Fixed Points: Classification

The flow near fixed point, x^* , for $\dot{x} = f(x)$

with $f(x^*) = 0$ is

$$\Delta x \approx e^{f'(x^*)t}$$

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What do points Δx distance from x^*
do as time gets large?

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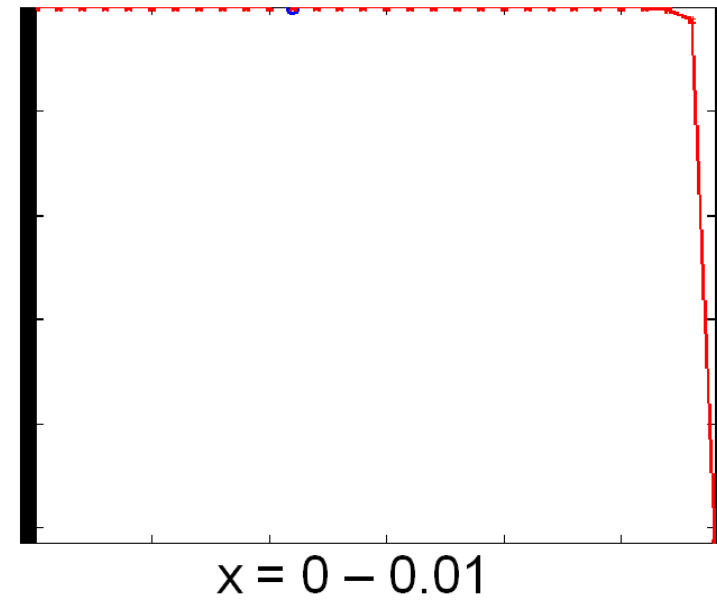
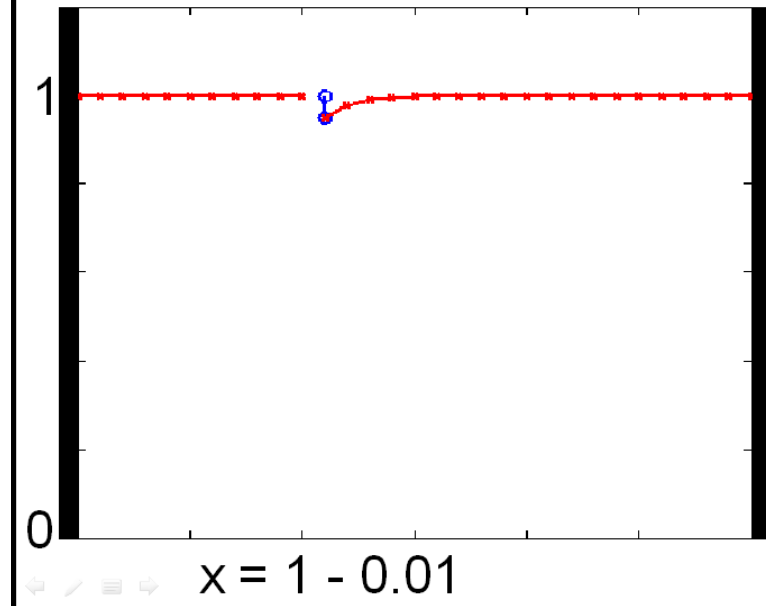
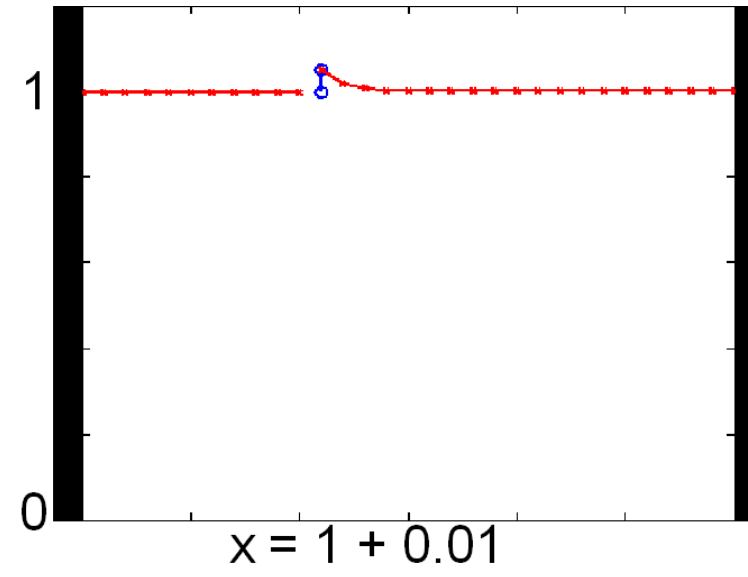
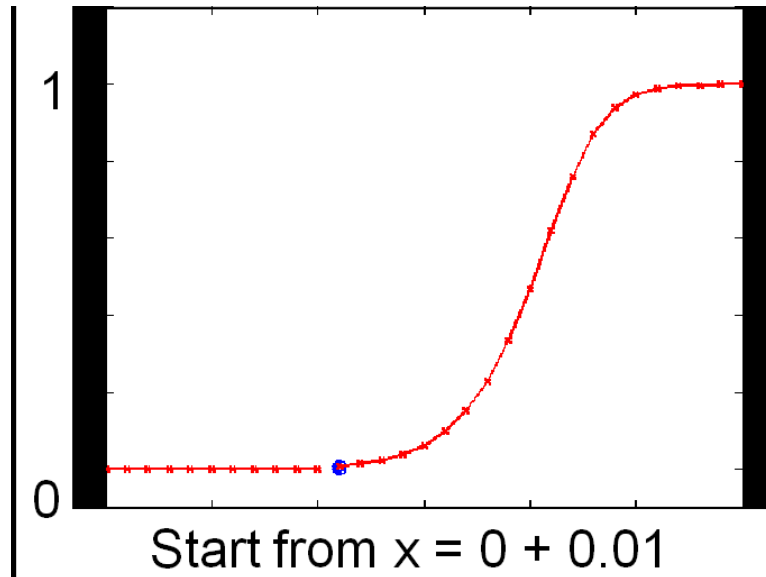
$$\lim_{t \rightarrow \infty} e^{f'(x^*)t} = \begin{cases} 0 & \text{if } f'(x^*) < 0 \\ \infty & \text{if } f'(x^*) > 0 \end{cases}$$

Fixed Points: Classification

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So for a 1D system the flow diverges away from the fixed point when the derivative of points near the fixed point are negative and converge on the fixed point if positive.

Fixed point at 0 is unstable.



Fixed point at 1 is stable.



Lyapunov Exponent

Aleksandr
Lyapunov,
Russia, turn of the
20th Century.

Fixed Points: Lyapunov Exponent

$$\lim_{t \rightarrow \infty} e^{f'(x^*)t} = \begin{cases} 0 & \text{if } f'(x^*) < 0 \\ \infty & \text{if } f'(x^*) > 0 \end{cases}$$

As we just showed the time evolution close to a fixed point x^* is generally exponential:

$$\Delta x = e^{\lambda t}, \text{ where } \lambda = f'(x^*).$$

Lyapunov Exponent

Fixed Points: Lyapunov Exponent

A negative Lyapunov Exponent \implies
the flow moves exponentially towards
the fixed point.

A positive Lyapunov Exponent \implies
the flow moves exponentially away from
the fixed point.

Fixed Points: Discrete Maps

$$x(t + 1) = f(x(t))$$

$$x(t + 1) = f(x(t)) \approx f(x^*) + g'(x^*)(x(t) - x^*)$$

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For discrete maps the fixed point is at:

$$f(x^*) = x^*$$

Fixed Points: Discrete Maps

Why?

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Fixed Points: Discrete Maps

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Fixed Points: Discrete Maps

$$\Delta x(t+1) = x(t+1) - x^*$$

Fixed Points: Discrete Maps

$$\Delta x(t+1) = x(t+1) - x^* = f'(x^*)\Delta x(t)$$

Fixed Points: Discrete Maps

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Fixed Points: Discrete Maps Solve for λ

$$\Delta x(t+1) = x(t+1) - x^* = f'(x^*)\Delta x(t)$$

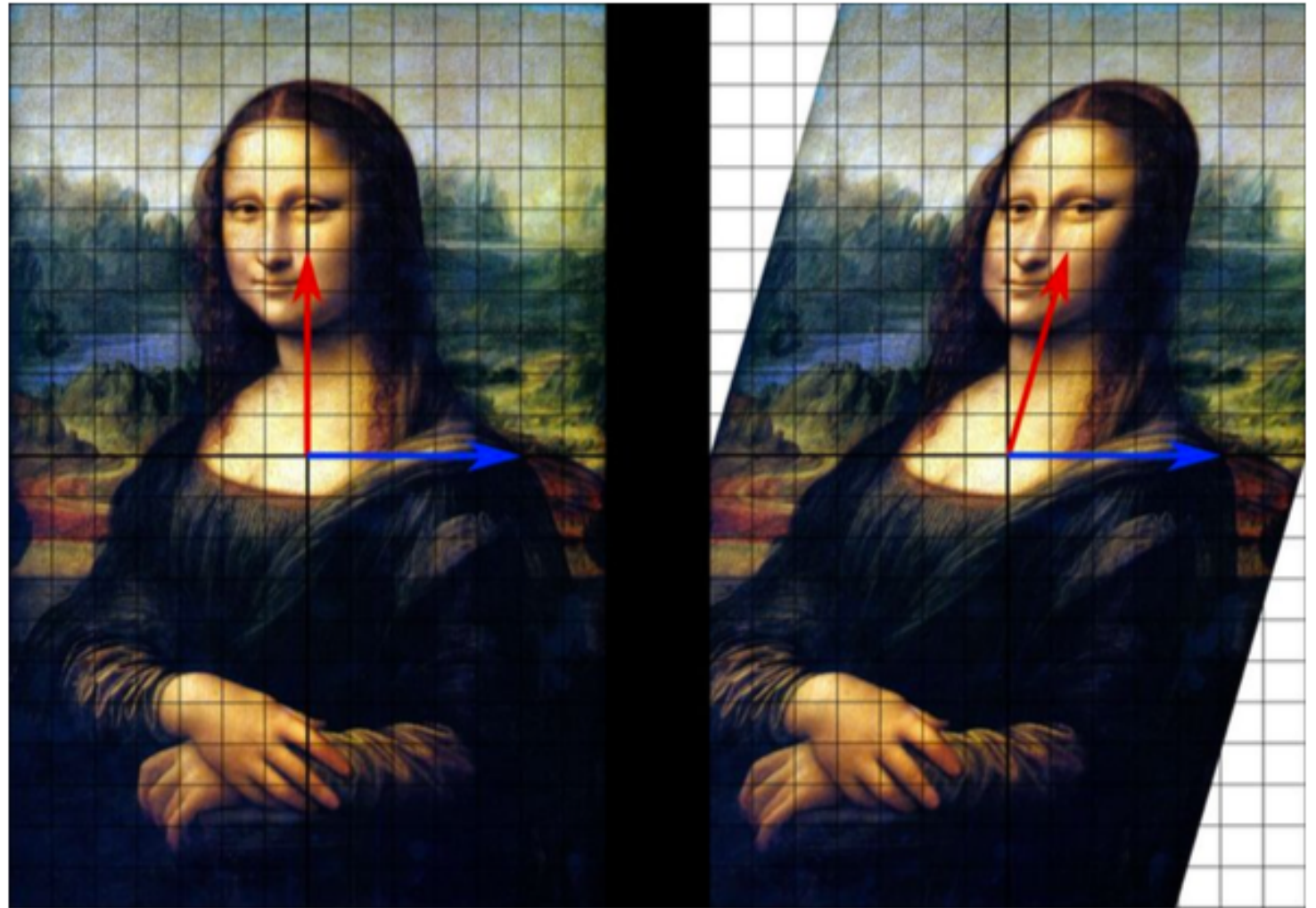
$$\Delta x(t+1) = x(t+1) - x^* = f'(x^*)\Delta x(t)$$

$$\Delta x(t) = e^{\lambda t}$$

The Lyapunov exponent for maps is:

$$\lambda = \ln |f'(x^*)| = \begin{cases} < 0 & \text{if } |f'(x^*)| < 1 \\ > 0 & \text{if } |f'(x^*)| > 1 \end{cases}$$

Fixed Points: Multidimensional Systems



The blue line is an eigenvector.

Wikipedia, 2017

It doesn't change as we transform the image.

Fixed Point Analysis: Multidimensional Systems

Consider the 2D system:

$$\dot{x} = y$$

$$\dot{y} = 2x + y$$

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Fixed Point Analysis: Multidimensional Systems

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In matrix form:

$$\begin{bmatrix} 0x + 1y \\ 2x + 1y \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Fixed Point Analysis: Multidimensional Systems

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In matrix form:

$$\begin{bmatrix} 0x + 1y \\ 2x + 1y \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 & 1 \\ 2 & 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

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The eigenvalues describe the flow nearby

Fixed Point Analysis: Multidimensional Systems

An enormous shortcut . . .

Fixed points are where:

In matrix form:

$$y = 0$$

$$2x + y = 0$$

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The eigenvalues describe the flow nearby

```
>> M = [0,1;2,1]
```

```
M =
```

```
    0    1  
    2    1
```

```
>> eig(M)
```

```
ans =
```

```
   -1  
    2
```

Fixed Point Analysis: Multidimensional Systems

The eigenvalues describe the flow nearby...

...if the system is LINEAR (or nearly so).

If non-linear we need to take the eigenvalues of the **Jacobian** matrix (see Lecture 7).

Jacobian

Given a set of equations:

$$y_1 = f_1(x_1, x_2, \dots, x_n)$$

$$y_2 = f_2(x_1, x_2, \dots, x_n)$$

⋮

$$y_m = f_m(x_1, x_2, \dots, x_n)$$

Jacobian

The Jacobian is:

The partial derivatives for each equation and in each direction.

(Calc III)

$$\begin{bmatrix} \frac{\partial y_1}{\partial x_1} & \frac{\partial y_1}{\partial x_2} & \cdots & \frac{\partial y_1}{\partial x_n} \\ \frac{\partial y_2}{\partial x_1} & \frac{\partial y_2}{\partial x_2} & \cdots & \frac{\partial y_2}{\partial x_n} \\ \vdots & \vdots & \ddots & \vdots \\ \frac{\partial y_m}{\partial x_1} & \frac{\partial y_m}{\partial x_2} & \cdots & \frac{\partial y_m}{\partial x_n} \end{bmatrix}$$

Example: Holmes Map

$$x_{t+1} = y_t$$

$$y_{t+1} = -bx_t + ay_t - y_t^3$$

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$$y_{t+1} = -bx_t + ay_t - y_t^3$$

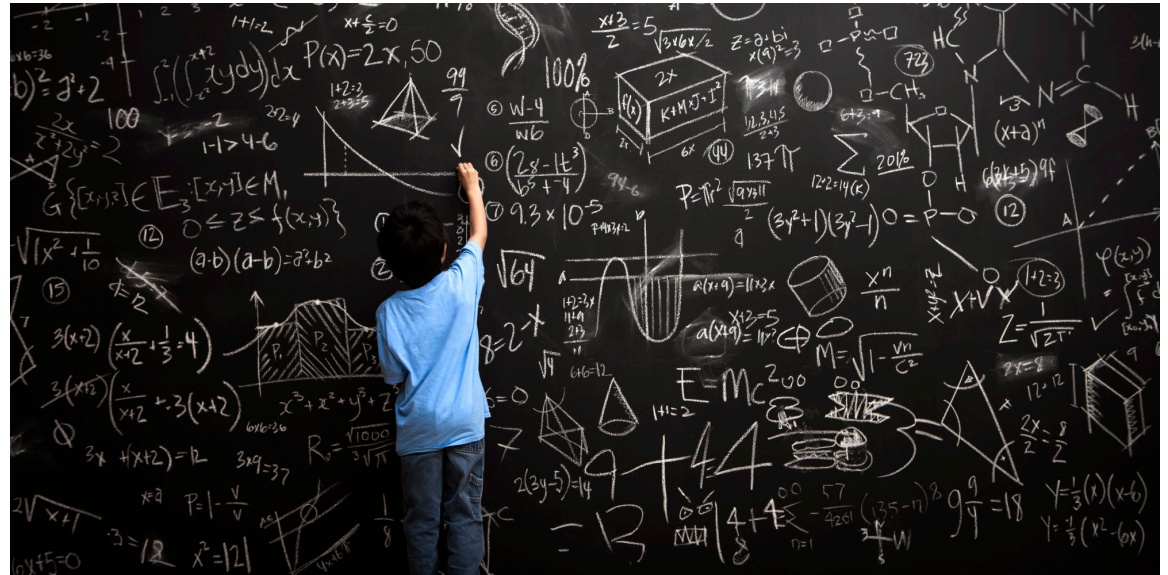
$$J = \begin{bmatrix} \frac{\partial y_t}{\partial x_t} & \frac{\partial y_t}{\partial y_t} \\ \frac{\partial(-bx_t + ay_t - y_t^3)}{\partial x_t} & \frac{\partial(-bx_t + ay_t - y_t^3)}{\partial y_t} \end{bmatrix}$$

The image shows the Jacobian matrix J for the Holmes map. The top-left element is $\frac{\partial y_t}{\partial x_t}$, the top-right is $\frac{\partial y_t}{\partial y_t}$, the bottom-left is $\frac{\partial(-bx_t + ay_t - y_t^3)}{\partial x_t}$, and the bottom-right is $\frac{\partial(-bx_t + ay_t - y_t^3)}{\partial y_t}$. Blue arrows point from the labels x_{t+1} to the y_t terms in the top row of the matrix.

Example: Holmes Map

$$J = \begin{bmatrix} \frac{\partial y_t}{\partial x_t} & \frac{\partial y_t}{\partial y_t} \\ \frac{\partial(-bx_t + ay_t - y_t^3)}{\partial x_t} & \frac{\partial(-bx_t + ay_t - y_t^3)}{\partial y_t} \end{bmatrix}$$

Let's Calculate it...



Example: Holmes Map

$$J = \begin{bmatrix} \frac{\partial y_t}{\partial x_t} & \frac{\partial y_t}{\partial y_t} \\ \frac{\partial(-bx_t + ay_t - y_t^3)}{\partial x_t} & \frac{\partial(-bx_t + ay_t - y_t^3)}{\partial y_t} \end{bmatrix}$$
$$= \begin{bmatrix} 0 & 1 \\ -b & a - 3y^2 \end{bmatrix}$$

In Matlab:

```
>> syms x y a b
```

```
>> HolmesMap = [y; -b*x+a*y-y^3]
```

```
HolmesMap = y - y^3 + a*y - b*x
```

```
>> HolmesMapJ = jacobian(HolmesMap, [x,y])
```

```
HolmesMapJ = [0, 1]  
             [-b, -3*y^2 + a]
```

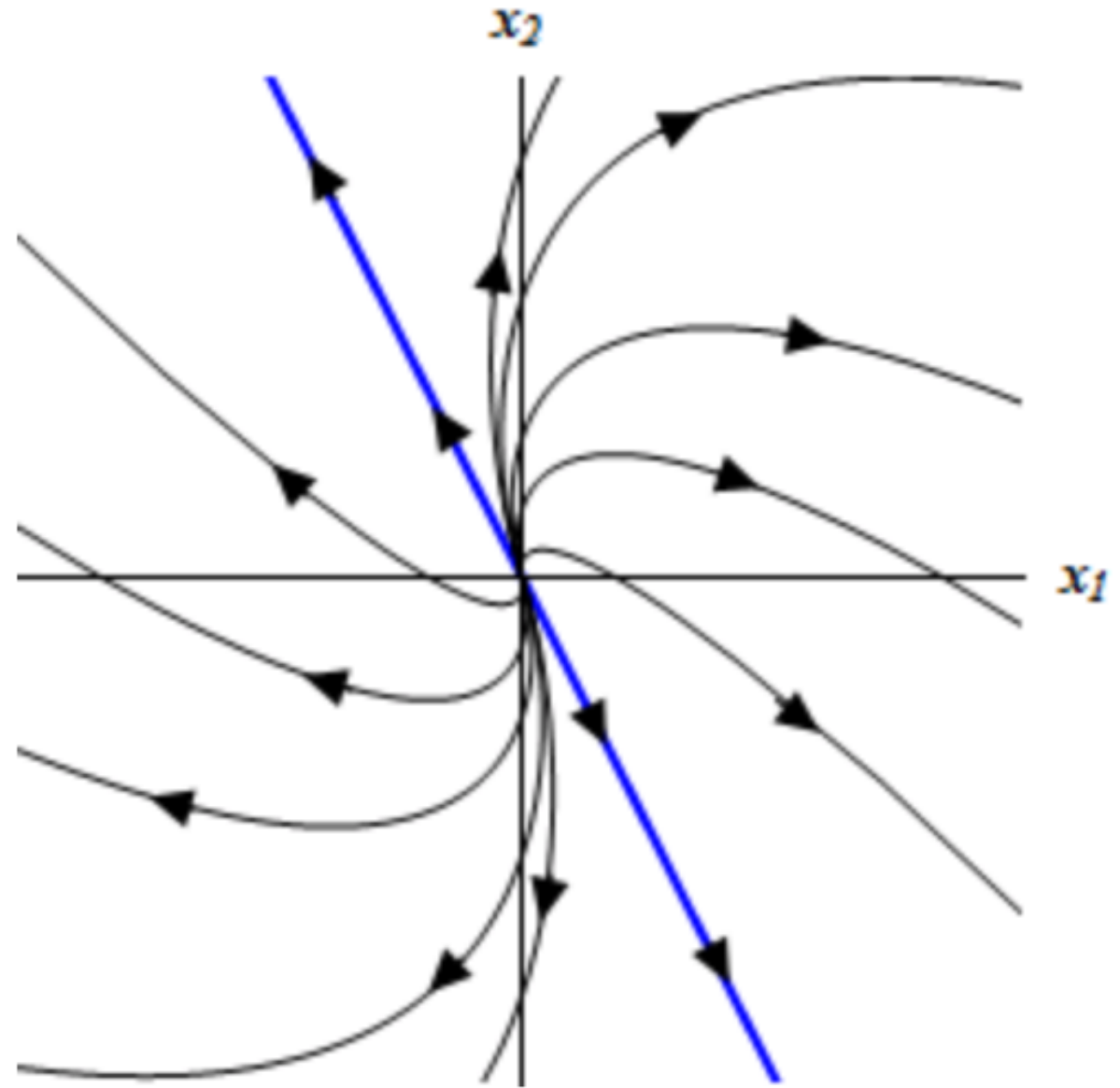
```
>> eig(HolmesMapJ)
```

```
ans = a/2 - (a^2 - 6*a*y^2 + 9*y^4 - 4*b)^(1/2)/2 - (3*y^2)/2  
      a/2 + (a^2 - 6*a*y^2 + 9*y^4 - 4*b)^(1/2)/2 - (3*y^2)/2
```

Classifying Fixed Points (2D Systems)

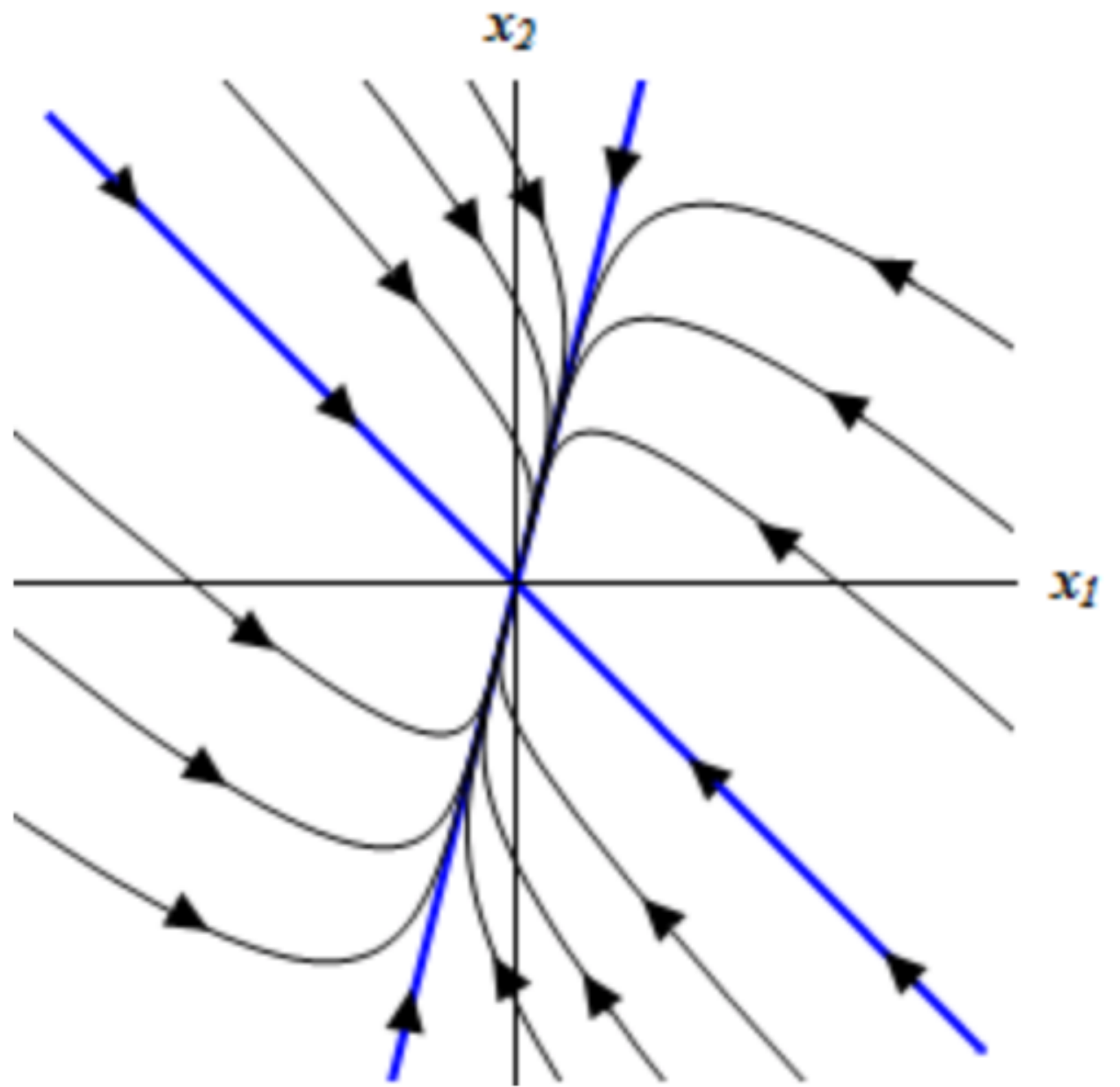
Eigenvalues	Stability	Name
Real and positive	Unstable	Source
Real and negative	Stable	Sink
Real mixed signs	Unstable	Saddle point
Complex with positive real part	Unstable	Spiral Source
Complex with negative real part	Stable	Spiral Sink
Imaginary	Unstable	Center

Classifying
Fixed-Points
(2D)



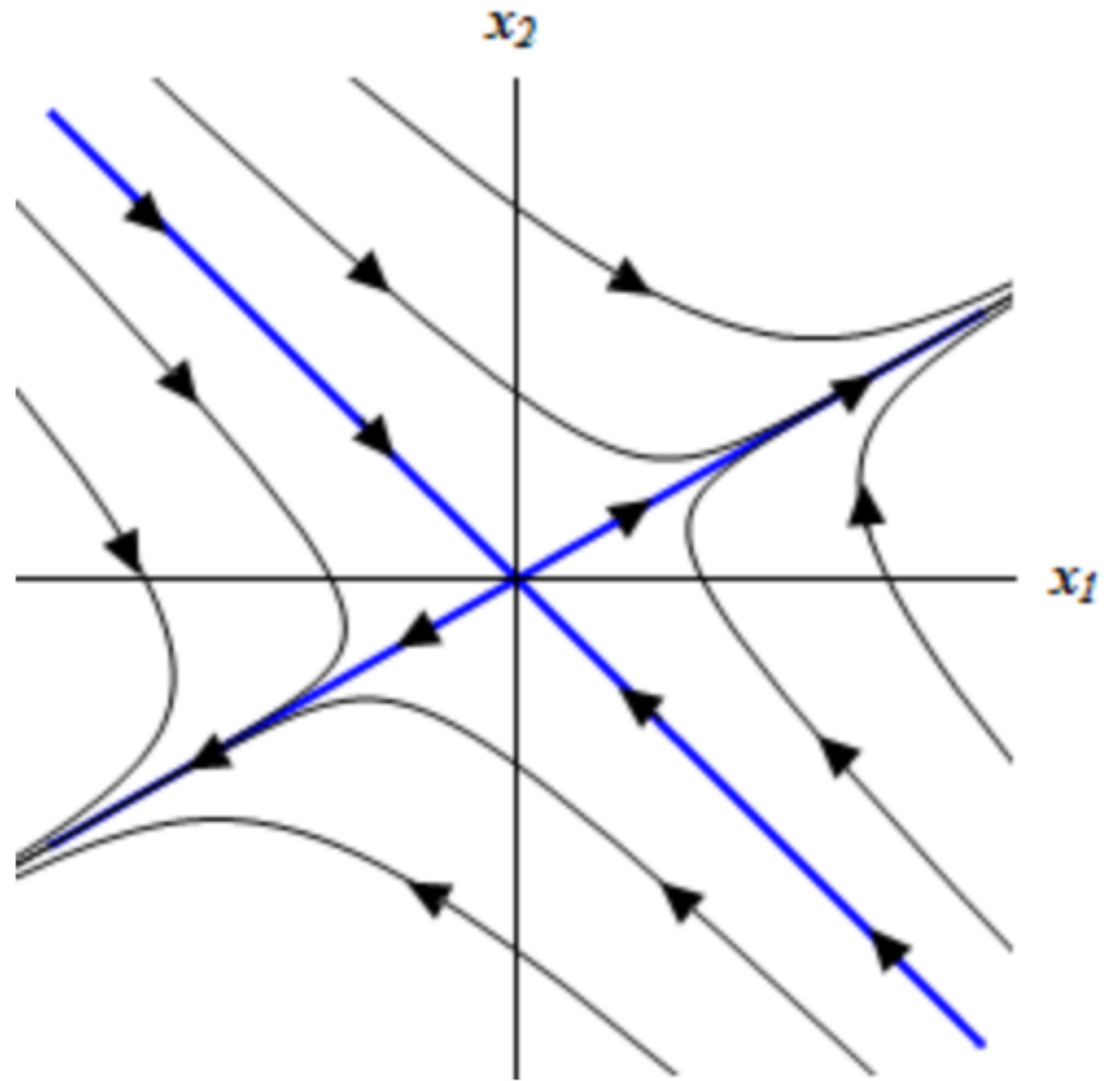
Unstable fixed-point (source)

Classifying
Fixed-Points
(2D)



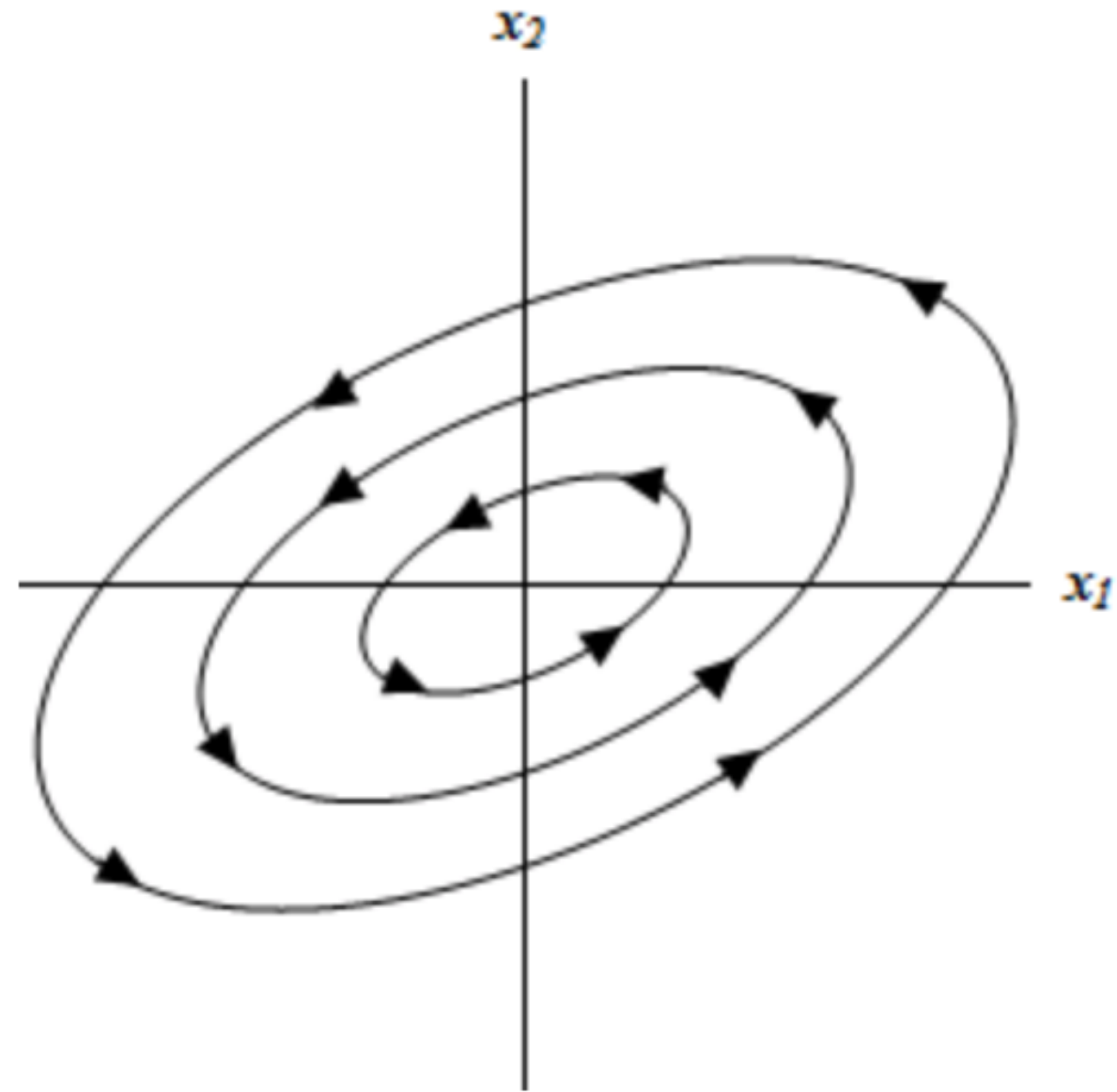
Stable fixed-point (sink)

Classifying
Fixed-Points
(2D)



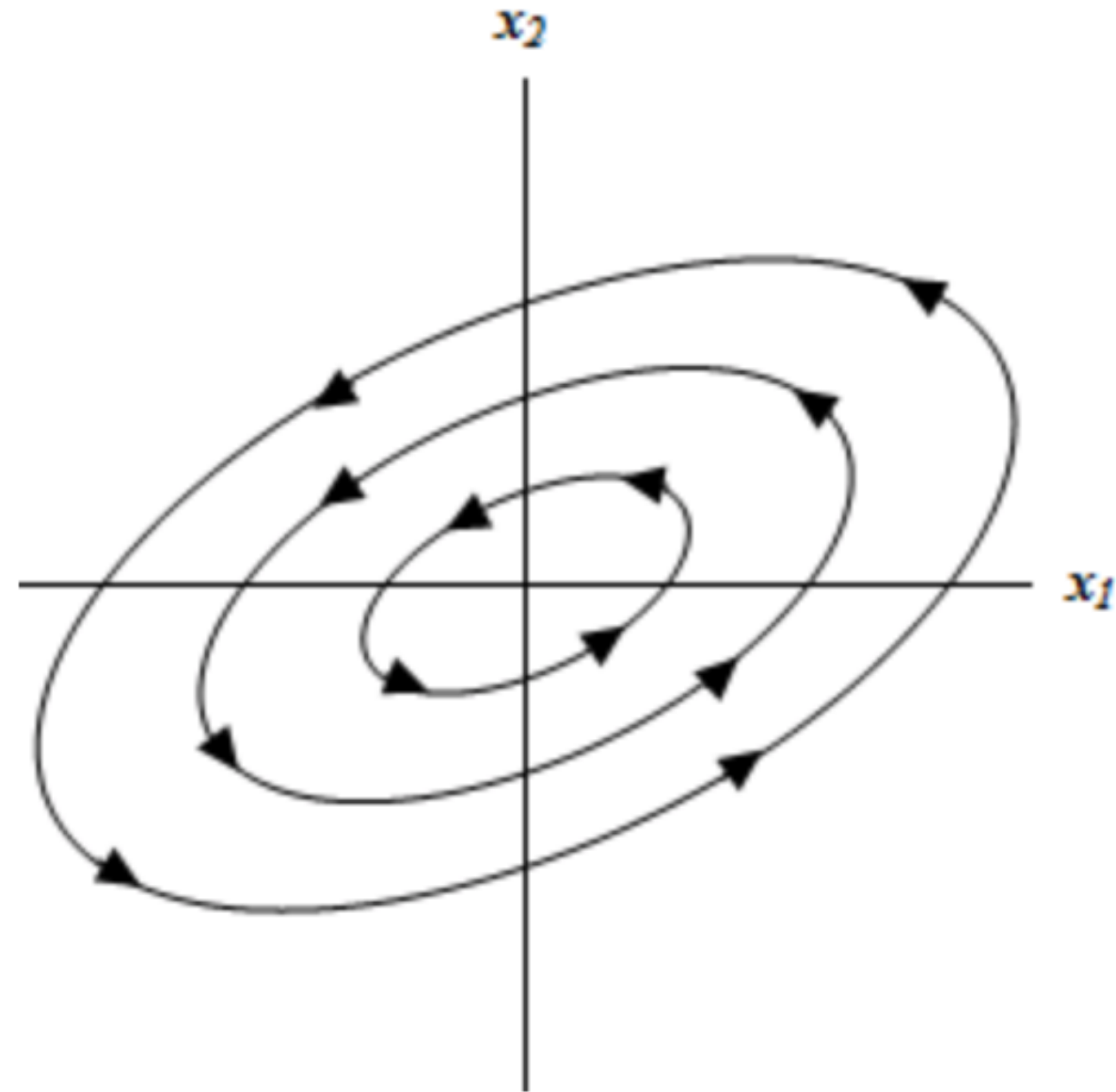
Saddle point

Classifying
Fixed-Points
(2D)



Center

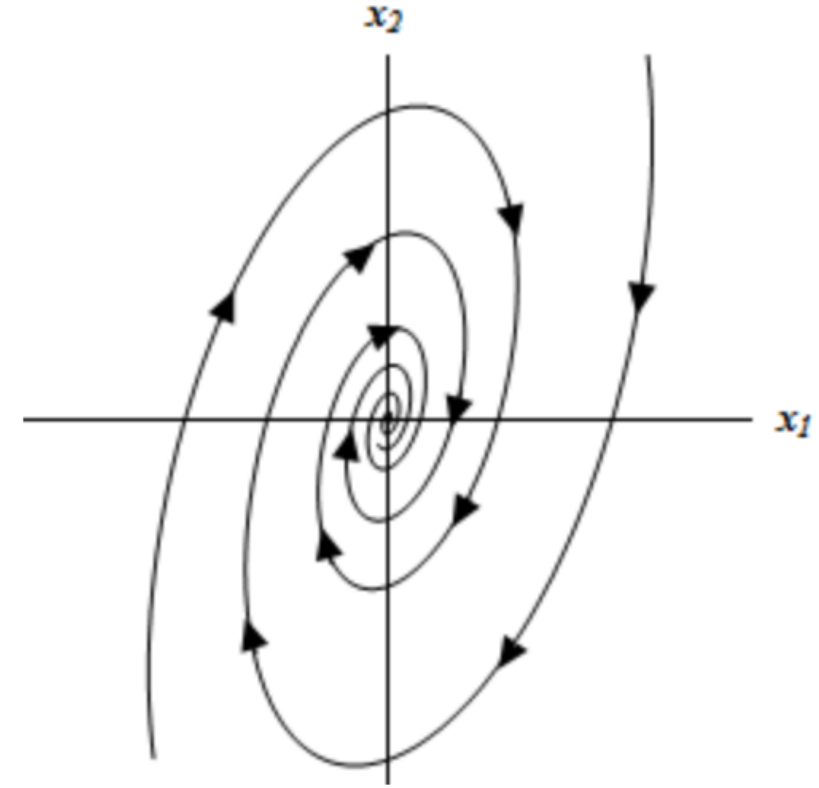
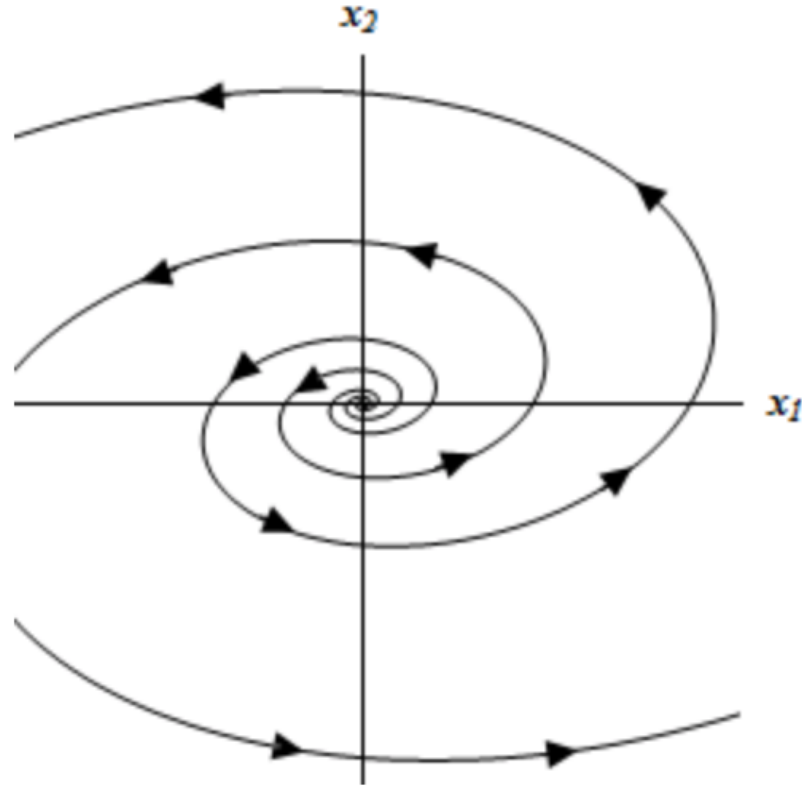
Classifying Fixed-Points (2D)



Center

Recall the Lotke-Volterra Model

Classifying
Fixed-Points
(2D)



Spirals (stable and unstable)

Lecture 5

Ergodicity

Ergodicity

A dynamical system in which trajectories come arbitrarily close to any point in the phase space no matter the initial conditions.

This implies that the time average is equal to the spacial average.

Basic Concepts: Ergodicity

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is equal to the spacial average.

So what?

Basic Concepts: Ergodicity

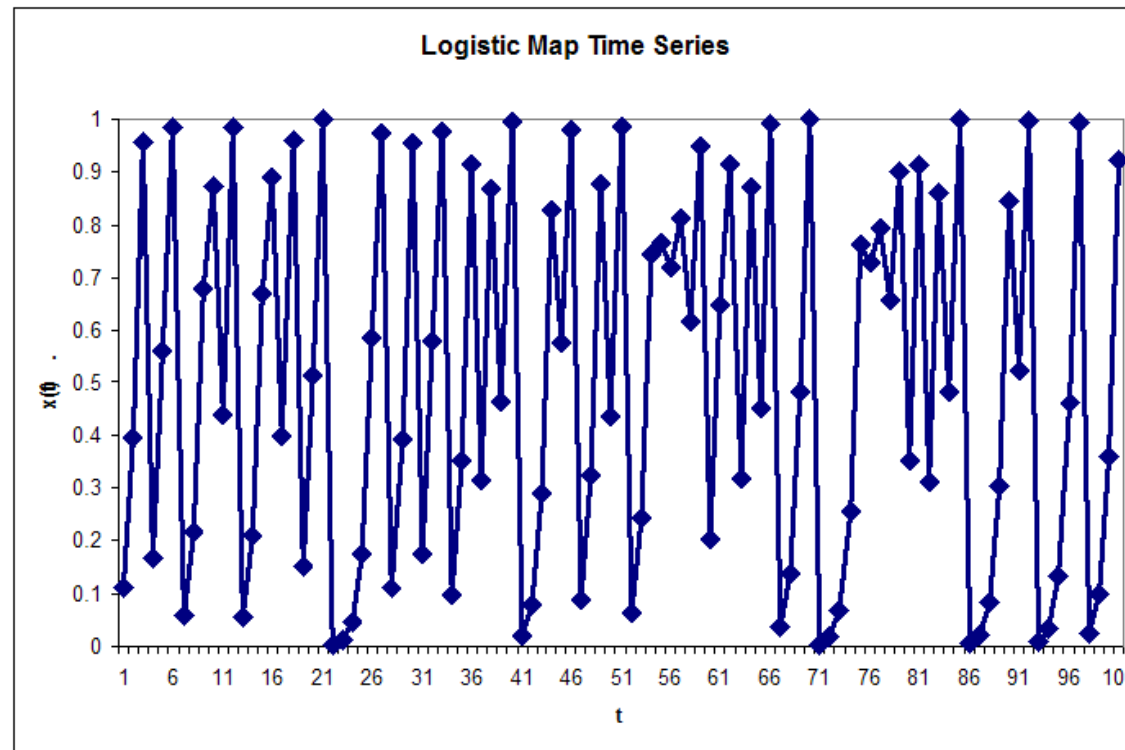
This implies that the time average is equal to the spacial average.

So what?

If a system is ergodic we can make good long-term average predictions even when the system is chaotic.

Recall the chaos of the logistic map

$$x_{t+1} = rx_t(1 - x_t)$$



The logistic map turns out to be ergodic

Ergodicity: Logistic Map

Histogram of 4000 randomly chosen initial conditions.

Histogram of time steps > 2000 for a single initial condition.

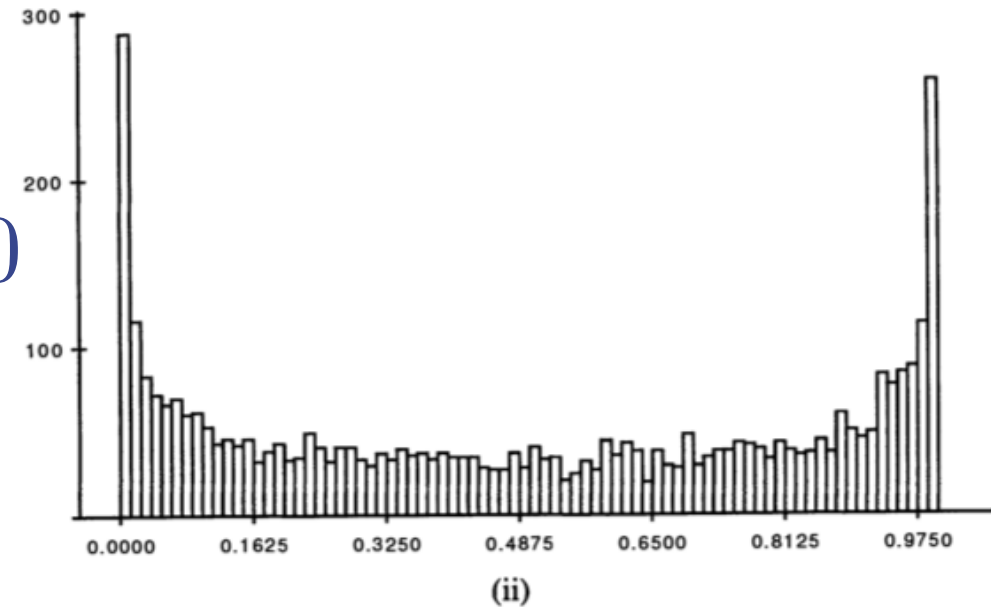
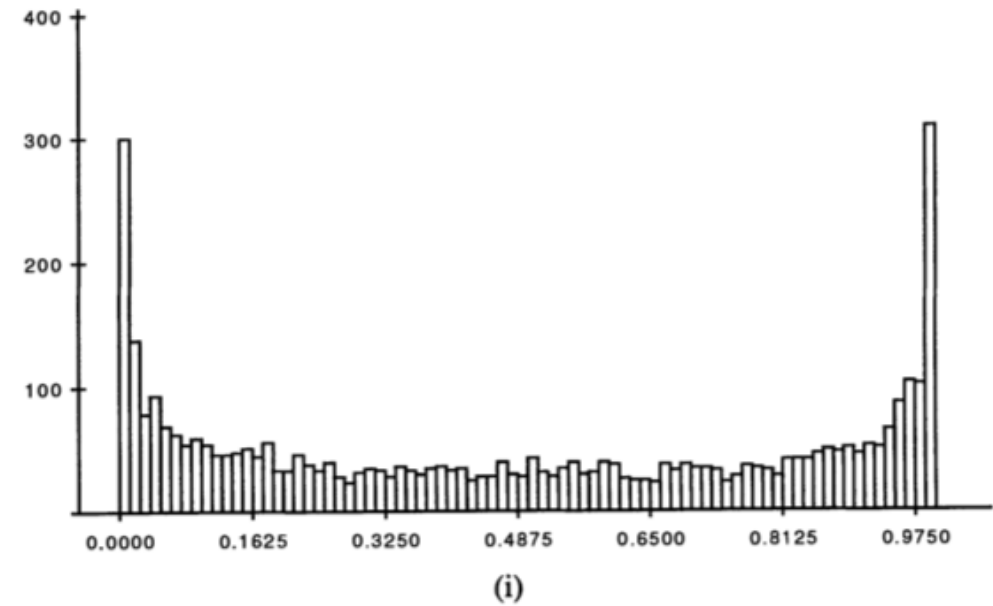
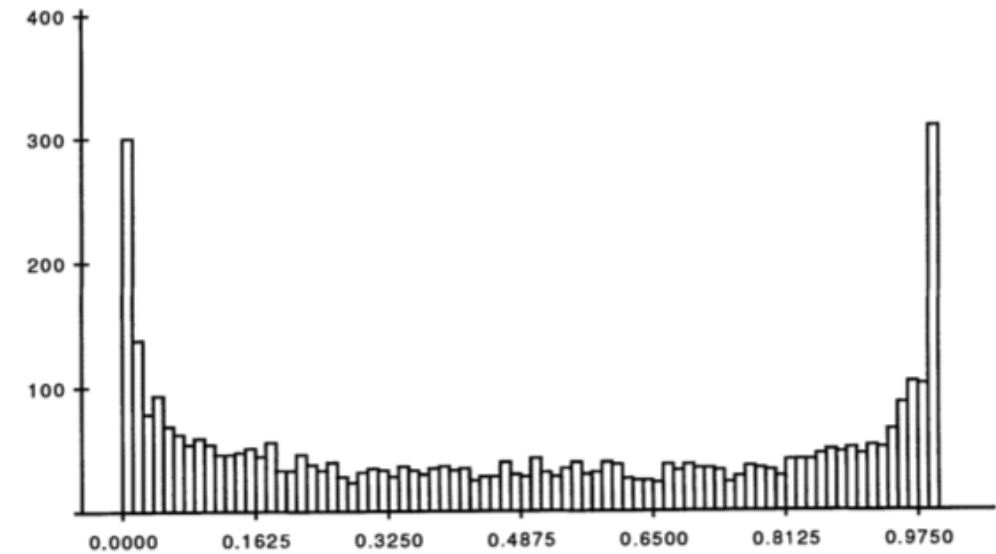


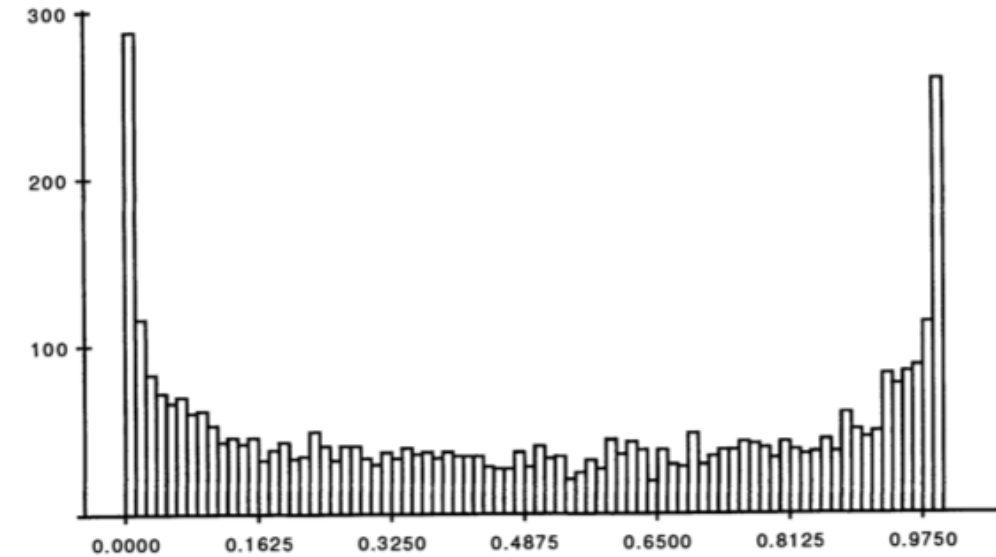
FIG. 8. Example of ergodic behavior: Logistic map, $a = 4.0$. (i) Histogram of 4000 iterates of $x_0 = .20005$. (ii) Histogram of the logistic map at time 2000 for 4000 x_0 's in $[.10005, .30005]$.

Ergodicity: Logistic Map

Because the Logistic map is ergodic we can predict its average behaviour.



(i)



(ii)

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Ergodicity: Logistic Map

Ergodic systems never become trapped in a particular region of phase space. They eventually roam everywhere.

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So can an ergodic system have attractors?

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So can an ergodic system have attractors?

Is the logistic map for $r=1$ ergodic?